

Density of states in 1-dimension

Let's consider a chain of particles in 1D.
 $N+1$ particles, separated by a distance a , in
such a way that the chain has length L

→ The particles at $s=0$ and $s=N$ at the extremes
of the chain/line are fixed.

→ Each normal mode of polarization p has the
form of a standing wave

If the amplitude of the oscillation of the s th
particle is u_s

$$u_s = u e^{-i\omega_{k,p}t} \sin(ska)$$

$\omega_{k,p}$ is related through k to the dispersion
relation.

To force the amplitude at $s=0$ and $s=N$ to be
zero we choose

$$k = \frac{n\pi}{L} \quad n \in \mathbb{Z}$$

$$\sin(ska) = \sin\left(\frac{sn\pi a}{L}\right) \quad \text{at } n=N \quad \sin\left(s\pi \frac{Na}{L}\right)$$

$s=0$ $s=N$

$L = Na$

If $n=N$, we have $u_s = 0$ for all particles

$$k = \frac{\pi}{L}, \frac{2\pi}{L}, \frac{3\pi}{L}, \dots, \frac{(N-1)\pi}{L}$$

A set of $(N-1)$ values of k , which is equal to the number of particles that are allowed to move.

In a one-dimensional line of particles there is one mode for every interval $\Delta k = \frac{\pi}{L}$

The number of modes per unit range k

$$\text{is } \frac{1}{\Delta k} = \frac{L}{\pi} \quad \text{for } k \leq \pi/a \quad k = \frac{n\pi}{L} = \frac{\pi}{a}$$
$$0 \quad \text{for } k > \pi/a$$

Other ways of counting are equally valid.

We can also consider a case where the line is not bounded, but we require the solutions to be periodic over the distance L .

$$u_s(sa) = u_s(sa + L)$$

This method is referred to as the method of periodic boundary conditions

Now we have a travelling wave

$$u_s = u e^{i(ska - \omega_{kp}t)}$$

For this wave to be periodic we need

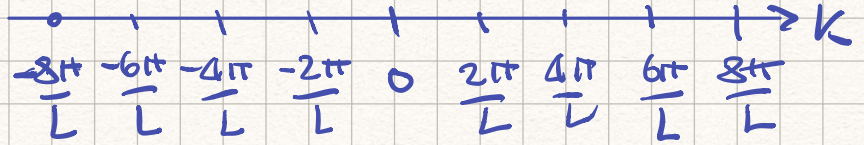
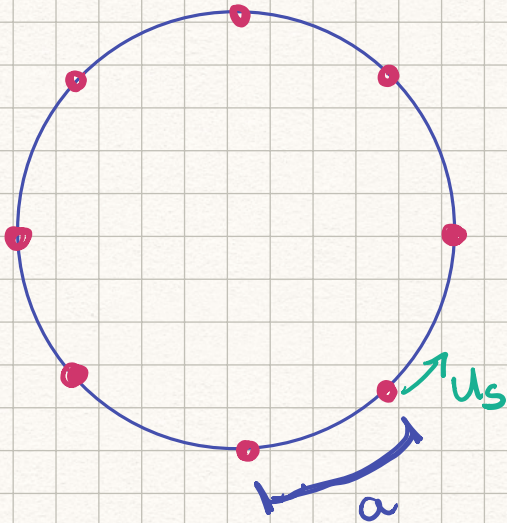
$$e^{iks_a} = e^{ik(s_a + L)} = e^{iks_a} e^{ikL} \Rightarrow 1 = e^{ikL}$$

$$kL = \pm 2n\pi \quad \Rightarrow \quad k = \pm \frac{2n\pi}{L}$$

The valid k -values are $\pm \frac{2\pi}{L}, \pm \frac{4\pi}{L}, \pm \frac{6\pi}{L}, \dots$

$$N=8$$

$k=0$: uniform mode



There are 8 modes of oscillation

The amplitude of oscillation

$$\perp, e^{\pm i\pi s/4}$$

$$L=8a$$

$$u_s = u e^{i(ska - \omega_{kp}t)}$$

$$k = \pm \frac{2\pi}{L} = \pm \frac{2\pi}{8a}$$

$$= \pm \frac{\pi}{4a}$$

$$k=0$$

$$u_s = \perp u e^{-i\omega t}$$

$$k = \pm \frac{2\pi}{L}$$

$$u_s = e^{\pm i\pi s/4} u e^{-i\omega t}$$

$$e^{i s \frac{\pi}{4a} a} = e^{i\pi s/4}$$

$$k = \pm \frac{4\pi}{L}$$

$$u_s = e^{\pm i\pi s/2} u e^{-i\omega t}$$

$$e^{i s \frac{6\pi}{8a} a} =$$

$$k = \pm \frac{6\pi}{L}$$

$$u_s = e^{\pm i3\pi s/4} u e^{-i\omega t}$$

$$e^{i s \frac{8\pi}{8a} a} = e^{i\pi s}$$

$$k = -\frac{8\pi}{L}$$

$$u_s = e^{-i\pi s} u e^{-i\omega t}$$

We have one solution (for k) per mobile atom but now we have \pm values of k

$$\Delta k = \frac{2\pi}{L}$$

$$-\frac{\pi}{a} < k < \frac{\pi}{a} : \frac{L}{2\pi} \left\{ \begin{array}{l} \text{density of } k \\ \text{values} \end{array} \right.$$

otherwise : 0

We need to obtain $D(\omega)$, the number of modes per unit frequency range for a given polarization.

In 1D:

$$D_1(k) d\omega = \frac{L}{\pi} \frac{dk}{d\omega} d\omega$$

$$= \frac{L}{\pi} \frac{1}{v_g} d\omega$$

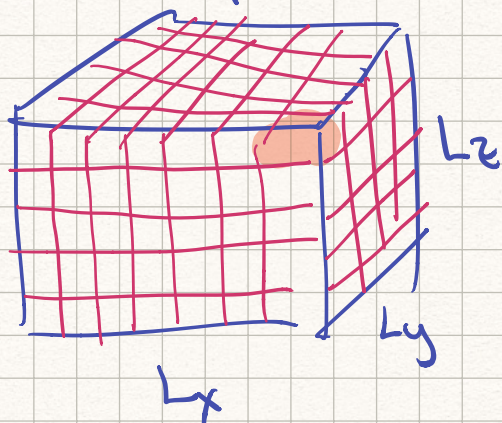
$$\frac{dk}{d\omega} = \frac{1}{d\omega/dk} = \frac{1}{\frac{d\omega}{dk}}$$

$$v_g = \frac{d\omega}{dk} \quad \text{Related to the dispersion relation}$$

Note that we have a singularity in $D_1(\omega)$ whenever the dispersion relation $\omega(k)$ is horizontal $\Rightarrow v_g = 0$

Density of states in three dimensions

Periodic bound. cond. method N^3 cells within a cube of side $L_x \cdot L_y \cdot L_z$



$$Na_x = L_x$$

$$Na_y = L_y$$

$$Na_z = L_z$$

$$a_x \cdot a_y \cdot a_z$$

$$\vec{k} = (k_x, k_y, k_z)$$

$$\vec{L} = (L_x, L_y, L_z)$$

$$\vec{r} = (x, y, z)$$

$$e^{i\vec{k} \cdot \vec{r}} = e^{i\vec{k} \cdot (\vec{r} + \vec{L})}$$

$$\vec{k} \cdot \vec{r} = (k_x x + k_y y + k_z z)$$

$$\vec{r} + \vec{L} = (x + L_x, y + L_y, z + L_z)$$

$$e^{ik_x x} e^{ik_y y} e^{ik_z z} = e^{ik_x (x + L_x)} e^{ik_y (y + L_y)} e^{ik_z (z + L_z)}$$

So that

$$k_x = 0; \pm \frac{2\pi}{L_x}; \pm \frac{4\pi}{L_x} \dots; -\frac{N\pi}{L_x}$$

$N = \text{even number}$

$$k_y = 0; \pm \frac{2\pi}{L_y}; \pm \frac{4\pi}{L_y} \dots; -\frac{N\pi}{L_y}$$

$$k_z = 0; \pm \frac{2\pi}{L_z}; \pm \frac{4\pi}{L_z} \dots; -\frac{N\pi}{L_z}$$

There's 1 \vec{k} -vector per unit volume $\frac{(2\pi)^3}{L_x L_y L_z}$

$$\frac{(2\pi)^3}{L_x L_y L_z} = \frac{2\pi}{L_x} \cdot \frac{2\pi}{L_y} \cdot \frac{2\pi}{L_z}$$

The density of \vec{k} -vectors is $\frac{L_x L_y L_z}{(2\pi)^3} = \frac{V}{(2\pi)^3}$

$V = L_x L_y L_z$: volume.

The total number of modes that have wavevector less than k is given by $\frac{V}{(2\pi)^3}$ multiplied by the volume of a sphere with radius k .

$$N = \frac{V}{(2\pi)^3} \frac{4\pi k^3}{3} = \frac{1}{3} \frac{V k^3}{2\pi^2}$$

Therefore, the density of states for each polarization is

$$\begin{aligned} \mathcal{D}_3(\omega) d\omega &= \frac{\partial N}{\partial \omega} d\omega = \frac{\partial N}{\partial k} \frac{\partial k}{\partial \omega} d\omega \\ &= \frac{V k^2}{2\pi^2} \frac{\partial k}{\partial \omega} d\omega \quad \omega(k) \end{aligned}$$