

Phonons - continued

$$\omega^2 = \frac{c(m_1 + m_2) \pm c \sqrt{m_1^2 + m_2^2 + 2m_1 m_2 \cos(ka)}}{m_1 m_2}$$

$$\begin{aligned} m_1^2 + m_2^2 + 2m_1 m_2 \cos(ka) + 2m_1 m_2 - 2m_1 m_2 &= \\ (m_1^2 + m_2^2 + 2m_1 m_2) + 2m_1 m_2 [\cos(ka) - 1] &= \\ (m_1 + m_2)^2 + 2m_1 m_2 [\cos(ka) - 1] &= \\ (m_1 + m_2)^2 \left\{ 1 + \frac{2m_1 m_2}{(m_1 + m_2)^2} [\cos(ka) - 1] \right\} \end{aligned}$$

In the long wavelength limit, $k = \frac{2\pi}{\lambda} \ll \frac{1}{a}$

$$\cos(ka) \approx 1 - \frac{k^2 a^2}{2}$$

$$= (m_1 + m_2)^2 \left\{ 1 + \frac{2m_1 m_2}{(m_1 + m_2)^2} \left(-\frac{k^2 a^2}{2} \right) \right\} = (m_1 + m_2)^2 \left(1 - \frac{m_1 m_2 k^2 a^2}{(m_1 + m_2)^2} \right)$$

$$\omega^2 = \frac{c(m_1 + m_2) \pm c(m_1 + m_2) \sqrt{1 - \chi}}{m_1 m_2} \quad \chi = \frac{m_1 m_2 k^2 a^2}{(m_1 + m_2)^2}$$

$$\sqrt{1 - \chi} \approx 1 - \frac{1}{2}\chi$$

$$\omega^2 = \frac{c(m_1 + m_2)}{m_1 m_2} \left[1 \pm \left(1 - \frac{1}{2}\chi \right) \right]$$

$$\begin{aligned} &= \left\{ \begin{aligned} \frac{c(m_1 + m_2)}{m_1 m_2} \left(2 - \frac{1}{2}\chi \right) &\approx \frac{2c(m_1 + m_2)}{m_1 m_2} = \omega_1^2 = 2c \left(\frac{1}{m_1} + \frac{1}{m_2} \right) \\ \frac{1}{2} \frac{c(m_1 + m_2)}{m_1 m_2} \chi &= \frac{c}{2} \frac{k^2 a^2}{m_1 + m_2} = \omega_2^2 \end{aligned} \right. \end{aligned}$$

The matrix that we had was

$$\begin{pmatrix} 0 \\ 0 \end{pmatrix} = \begin{pmatrix} m_1 \omega^2 - 2c & c(1 + e^{ika}) \\ c(e^{ika} + 1) & m_2 \omega^2 - 2c \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$$

In long-wavelength approx. $k \ll 1$ $e^{\pm ika} \approx 1 \pm ika$

$$\begin{pmatrix} 0 \\ 0 \end{pmatrix} \approx \begin{pmatrix} m_1 \omega^2 - 2c & c(2 - ka) \\ c(2 + ka) & m_2 \omega^2 - 2c \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$$

$$\text{For } \omega = \omega_1^2 = 2c \left(\frac{1}{m_1} + \frac{1}{m_2} \right)$$

$$m_1 2c \left(\frac{1}{m_1} + \frac{1}{m_2} \right) u - 2cu + c(2 - ka)v = 0$$

$$2cu \frac{m_1}{m_2} + c(2 - ka)v = 0$$

$$c(2 + ka)u + m_2 2c \left(\frac{1}{m_1} + \frac{1}{m_2} \right) v - 2cv = 0$$

$$c(2 + ka)u + 2c \frac{m_2}{m_1} v = 0$$

At $k=0$

$$2cu \frac{m_1}{m_2} + 2cv = 0$$

$$2cu + 2c \frac{m_2}{m_1} v = 0$$

$$u \frac{m_1}{m_2} = -v$$

$$\frac{u}{v} = -\frac{m_2}{m_1}$$

Oscillation with opposite phases

$$\text{For } \omega^2 = \omega_2^2 = \frac{c}{2} \frac{k^2 a^2}{m_1 + m_2}$$

$$\frac{m_1 c k^2 a^2}{2(m_1 + m_2)} u - 2cu + c(2 - ka)v = 0$$

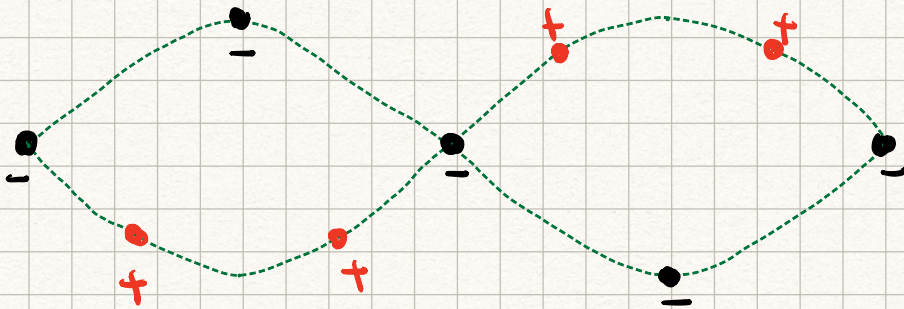
$$c(2+ka)u + \frac{m_2 c k^2 a^2}{2(m_1+m_2)}v - 2cv = 0$$

at $k=0$

$$-2cu + 2cv = 0 \rightarrow u = v$$

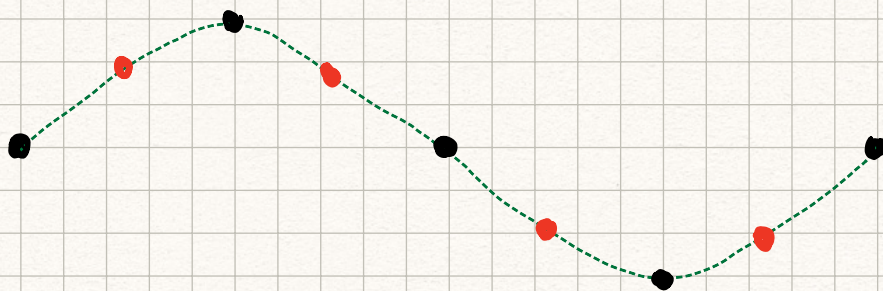
Oscillating in phase with the same amplitude

Optical mode: this is how a wave of light would make the charges move

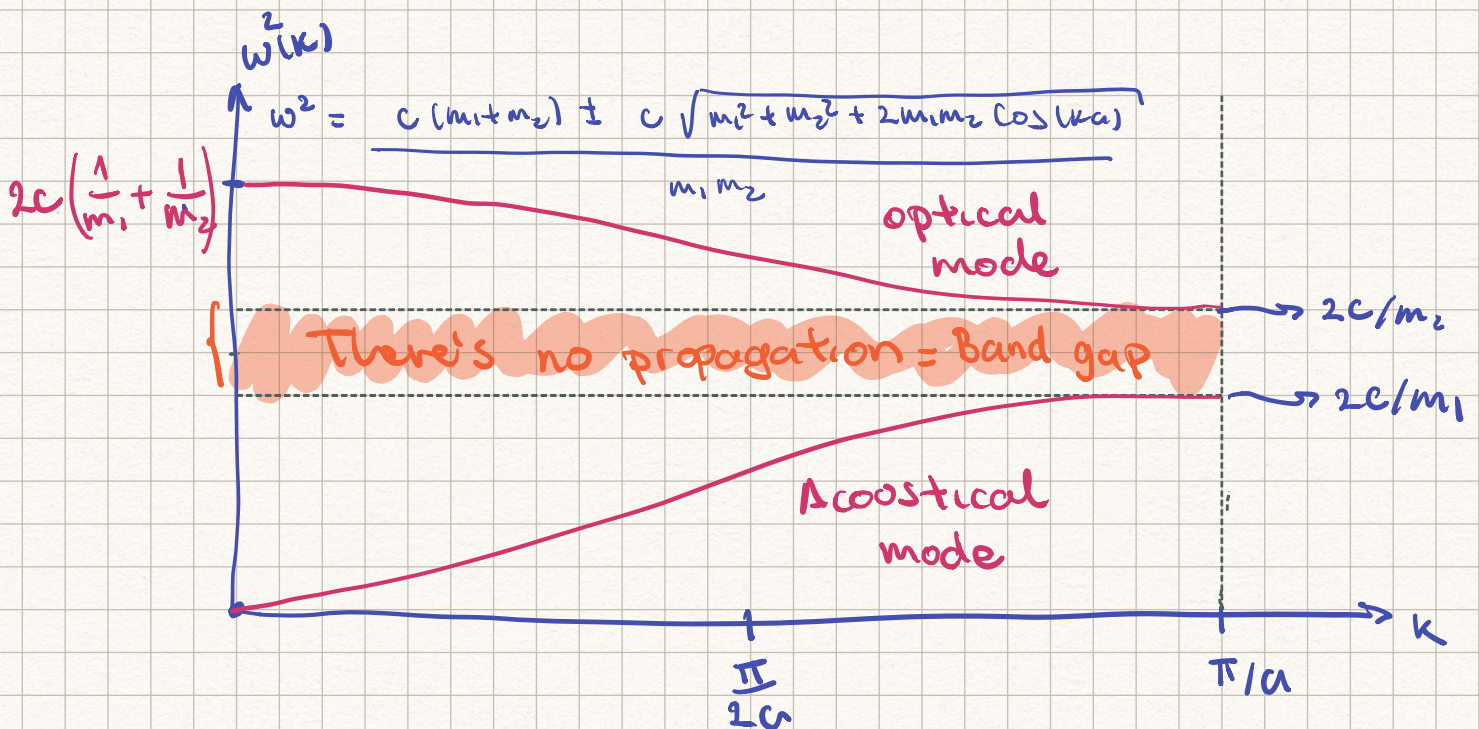


$$\frac{u}{v} = -\frac{m_1}{m_2}$$

Acoustic mode: this is how sound propagates



$$\frac{u}{v} = 1$$



$$k = \pi/a \quad \cos(\pi) = -1$$

$$\omega^2 = \frac{c(m_1 + m_2) \pm c \sqrt{m_1^2 + m_2^2 - 2m_1 m_2}}{m_1 m_2}$$

$$= \frac{c(m_1 + m_2) \pm c \sqrt{(m_1 - m_2)^2}}{m_1 m_2} \quad m_1 > m_2$$

$$= \frac{c(m_1 + m_2) \pm c(m_1 - m_2)}{m_1 m_2} = \begin{cases} 2c/m_2 \\ 2c/m_1 \end{cases}$$

$$k = 0 \quad \cos(0) = 1$$

$$\omega^2 = \frac{c(m_1 + m_2) \pm c \sqrt{m_1^2 + m_2^2 + 2m_1 m_2}}{m_1 m_2}$$

$$= \frac{c(m_1 + m_2) \pm c \sqrt{(m_1 + m_2)^2}}{m_1 m_2} \quad m_1 > m_2$$

$$= \frac{c(m_1 + m_2) \pm c(m_1 + m_2)}{m_1 m_2} = \begin{cases} 2 \frac{(m_1 + m_2)c}{m_1 m_2} = 2c \left(\frac{1}{m_1} + \frac{1}{m_2} \right) \\ 0 \end{cases}$$

There's a gap in the frequencies

$$\frac{2c}{m_1} < \omega^2 < \frac{2c}{m_2}$$

There are no wave-like solutions

In polyatomic lattices there's a frequency gap at the boundary of the Brillouin zone $k = \pm \pi/a$

Phonons are the ^{material} oscillations of the crystal