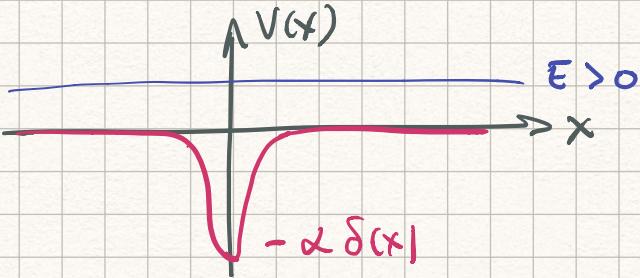


Scattering states of the δ -potential

$E > 0$: scattering states \Leftrightarrow not bounded

$$V(x) = -\infty \delta(x)$$



$$\left[\frac{p^2}{2m} + V(x) \right] \Psi(x) = E \Psi(x)$$

$$\left[-\frac{\hbar^2 \partial_x^2}{2m} - \infty \delta(x) \right] \Psi(x) = E \Psi(x) \Rightarrow \partial_x^2 \Psi + \frac{2mE}{\hbar^2} \delta(x) \Psi(x) = -\frac{2mE}{\hbar^2} \Psi(x)$$

$$= -k^2 \Psi(x)$$

$$k = \sqrt{\frac{2mE}{\hbar^2}} \quad E \in \mathbb{R}$$

Beyond the singularity ($\delta(x)$), except at $x=0$:

$$\partial_x^2 \Psi = -k^2 \Psi(x) \Rightarrow \Psi(x) = \begin{cases} A e^{ikx} + B e^{-ikx} & x < 0 \\ C e^{ikx} + D e^{-ikx} & x > 0 \end{cases}$$

$$\Psi(x=0) = \begin{cases} A+B \\ C+D \end{cases} \Rightarrow A+B = C+D$$

$$\int_{-\epsilon}^{\epsilon} \partial_x^2 \Psi dx + \frac{2m\delta}{\hbar^2} \int_{-\epsilon}^{\epsilon} \delta(x) \Psi(x) dx = -\frac{2mE}{\hbar^2} \int_{-\epsilon}^{\epsilon} \Psi(x) dx$$

Limit $\epsilon \rightarrow 0$

$$\Psi(\epsilon) - \Psi(-\epsilon) + \frac{2m\delta}{\hbar^2} \Psi(0) = 0 \quad (\times)$$

$$\Psi'(x) = \begin{cases} ik(A e^{ikx} - B e^{-ikx}) & x < 0 \\ ik(C e^{ikx} - D e^{-ikx}) & x > 0 \end{cases}$$

$$\Psi'(\epsilon) = ik(C e^{ik\epsilon} - D e^{-ik\epsilon}) = ik(C - D)$$

$$\Psi'(-\epsilon) = ik(A e^{-ik\epsilon} - B e^{ik\epsilon}) = ik(A - B)$$

$$\textcircled{X} \quad ik(C-D) - ik(A-B) = -\frac{2m\alpha}{\hbar^2} \psi(0)$$

$$ik(C-D-A+B) = -\frac{2m\alpha}{\hbar^2} (A+B)$$

$$C-D-A+B = -\frac{2m\alpha}{ik\hbar^2} (A+B)$$

$$C-D-A+B = 2i\beta(A+B)$$

$$C-D = A-B + 2i\beta(A+B)$$

$$C-D = A(1+2i\beta) - B(1-2i\beta)$$

$$A+B = C+D$$

$$\boxed{-\frac{1}{i} = i}$$

Check

$$B = \frac{m\alpha}{\hbar^2 k}$$

4 unknowns + K
we only have 2 eqs.!!!

Normalization doesn't help \Rightarrow not normalizable!

$$\underline{Ae^{ikx}} + \underline{B\bar{e}^{ikx}} \quad \text{Try it!}$$

Physical interpretation

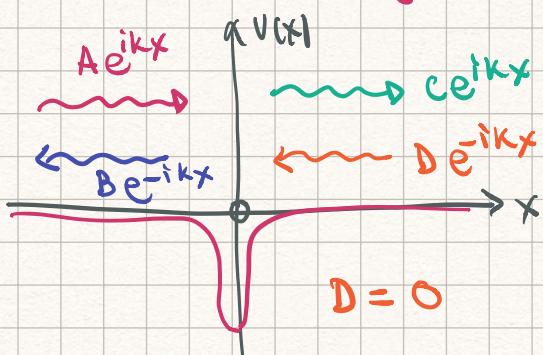
$$\psi(x,t) = \psi(x) e^{-iEt/\hbar}$$

$$= A e^{ik(x-Et/\hbar)}$$

wave that propagates to the right

$$+ B e^{-i(kx+Et/\hbar)}$$

wave that propagates to the left



\Rightarrow Scattering of particle that comes from the left \rightarrow

$\rightarrow A e^{ikx}$: Incoming

$\rightarrow C e^{ikx}$: Transmitted

$\rightarrow B \bar{e}^{ikx}$: Reflected

$$C = A(1+2i\beta) - B(1-2i\beta)$$

$$\underline{C = A+B}$$

$$A(1+2i\beta) - B(1-2i\beta) = A+B$$

$$A(1+2i\beta) - A = B + B(1-2i\beta)$$

$$2i\beta A = 2B(1-i\beta)$$

$$\boxed{B = \frac{i\beta}{1-i\beta} A}$$

$$C = A + B = A + \frac{iB}{1-iB} A = A \left[1 + \frac{iB}{1-iB} \right] = \boxed{A \frac{1}{1-iB} = C}$$

Scattering from the RIGHT, set $A=0$ and then

$$B = \frac{1}{1-iB} D \quad C = \frac{iB}{1-iB} D \quad \leftarrow D e^{-ikx}$$

$|N|^2$: density of prob x in space

Relative probability that the particle is

Reflected

$$R = \frac{|B|^2}{|A|^2} = \frac{|A|^2}{|A|^2} \left| \frac{iB}{1-iB} \right|^2$$

Reflection coefficient

$$\left| \frac{iB}{1-iB} \right|^2 = \left(\frac{iB}{1-iB} \right) \left(\frac{-iB}{1+iB} \right) = \frac{B^2}{1+B^2}$$

$$R = \frac{B^2}{1+B^2}$$

Transmitted

$$T = \frac{|C|^2}{|A|^2} = \frac{|A|^2}{|A|^2} \left| \frac{1}{1-iB} \right|^2$$

Transmission coefficient

$$\left| \frac{1}{1-iB} \right|^2 = \frac{1}{1+B^2}$$

$$T = \frac{1}{1+B^2}$$

of course $\boxed{R+T=1}$

$$B = \frac{m\alpha}{\hbar^2 k} \quad \text{and} \quad k = \frac{\sqrt{2mE}}{\hbar}$$

$$= \frac{m\alpha}{\hbar \sqrt{2mE}} \Rightarrow B^2 = \frac{m^2 \alpha^2}{\hbar^2 2mE} = \frac{m\alpha^2}{2E\hbar^2}$$

$$R = \frac{B^2}{1+B^2} = \frac{B^2}{B^2(1/\hbar^2 + 1)} = \frac{1}{1+1/\hbar^2}$$

$$1+\frac{1}{\hbar^2} = 1 + \frac{2E\hbar^2}{m\alpha^2}$$

$$T = \frac{1}{1+B^2}$$

$$R = \frac{1}{1 + 2E\hbar^2/(m\alpha^2)}$$

$$T = \frac{1}{1 + m\alpha^2/(2E\hbar^2)}$$

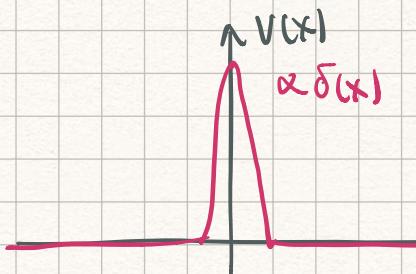
$$1+B^2 = 1 + \frac{m\alpha^2}{2E\hbar^2}$$

When $E \gg 1$: $R \rightarrow 0$
 $T \rightarrow 1$

Combination of stationary states \rightarrow wavepacket

\hookrightarrow range of energies

R and T are approximates for a particle with an energy in the vicinity of E.



$$\alpha \rightarrow -\alpha$$

- There's no bound state
-) check what happens when
 $\left. \begin{array}{l} \alpha \rightarrow -\alpha \\ \text{in the solutions of} \end{array} \right\}$ the previous lecture!

R and T are functions of $\alpha^2 \Rightarrow$ they remain the same

The particle is as likely to pass over a well, than through a barrier!!! $\xrightarrow{\text{transmitted}}$

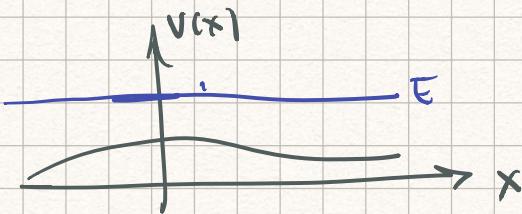
Classically

If $E \geq V_{\max}$:

$$R=0; T=1$$

$$E < V_{\max}$$

$$R=1; T=0$$



Quantum world

nonzero prob. go through $E < V_{\max}$
 \rightarrow Quantum tunneling !!!

$E > V_{\max}$: non zero prob. the particle will be reflected back.